Interpreting Neural Networks in the Context of Physical Phase Diagrams

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STRUCTURES CLUSTER OF EXCELLENCE



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Overview

Physics:

- > Phase Diagrams
- Standard Model

Machine Learning

- > Applying Neural Networks to Discover Phase Transitions
- Interpretation of Neural Networks

Phase Diagrams



Water





Atomic core

Magnet

Phase Diagrams



M Tc T Ferromagnet Paramagnet Magnet Water



Atomic core

Motivation

Theoretical Physics Goal:

Determine macroscopic phase diagrams from a microscopic description

- > Determine the existence of phases
- Pin down the phase transition
- Find the dominant characteristics of phases

Standard Model of Elementary Particles + Gravity



Standard Model of Elementary Particles + Gravity



Standard Model of Elementary Particles + Gravity



Standard Model of Elementary Particles + Gravity



Standard Model of Elementary Particles + Gravity



Experimentally verified paricle content + hypothetical graviton

Standard Model of Elementary Particles + Gravity



- We can only see part of the Standard Model, without
 - Heavy particles
 - Frozen Particles
 - Short Ranged Force Particles
- > Visible Matter consists of up- / down- quarks and electrons



> Ultimate Goal of the Talk: NN reveals nature of confinement PT

Standard Model of Elementary Particles + Gravity



Invitation: Phase transitions from microscopic physics







Supervised Learning 2d Ising Model

- Data: Monte Carlo samples
- Training at well known points in phase diagram
- Labels: Phase



- Testing in interval containing phase transition
- > Estimate within 1% of exact value $T_c = \frac{2}{\ln(1+\sqrt{2})}$



Machine Learning Phases of Matter

 Starting in 2016: Rush to calculate physical phase diagrams using Neural Networks



Machine learning phases of matter

Juan Carrasquilla¹⁴ and Roger G. Melko¹³

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Learning phase transitions by confusion

Evert P.L. van Nieuwenburg*, Ye-Hua Liu, and Sebastian D. Huber Institute for Theoretical Physics, ETH Zurich, 8093 Zürich, Switzerland

Classifying phases of matter is a central problem in physics. For quantum mechanical systems, this task can be daunting owing to the exponentially large Hilbert space. Thanks to the available computing power and access to ever larger data sets, classification problems are now routinely solved using machine learning techniques. Here, we propose to use a neural network based approach to find phase transitions depending on the performance of the neural network after training it with deliberately incorrectly labelled data. We demonstrate the success of this method on the topological phase transition in the Kitaev chain, the thermal phase transition in the classical Ising model, and the many-body-localization transition in a disordered quantum spin chain. Our method does not

Machine learning quantum phases of matter beyond the fermion sign problem

Peter Broecker,¹ Juan Carrasquilla,² Roger G. Melko,^{2,3} and Simon Trebst¹

¹Institute for Theoretical Physics, University of Cologne, 50937 Cologne, Germany ²Perimeter Institute for Theoretical Physics, Waterloo, Ontario N2L 2Y5, Canada ³Department of Physics and Astronomy, University of Waterloo, Ontario, N2L 3G1, Canada (Dated: August 30, 2016)

State-of-the-art machine learning techniques promise to become a powerful tool in statistical mechanics via their capacity to distinguish different phases of matter in an automated way. Here we demonstrate that convolutional neural networks (CNN) can be optimized for quantum many-fermion systems such that they correctly identify and locate quantum phase transitions in such systems. Using auxiliary-field quantum Monte Carlo (QMC) simulations to sample the many-fermion system, we show that the Green's function (but not the auxiliary field) holds sufficient information to allow for the distinction of different fermionic phases via a CNN. We demonstrate that this QMC + machine learning approach works even for systems exhibiting a severe fermion sign problem where conventional approaches to extract information from the Green's function, e.g. in the form of equal-time correlation functions, fail. We expect that this conscience works of hereneytical working mechanics that the severe of the strategian mechanics and the strategian technics approaches to extract information from the Green's function, e.g. in the form

niques to circumvent the fermion sign p problems in statistical physics. Matthew J. S. Beach,^{*} Anna Golubeva, and Roger G. Melko Department of Physics and Astronomy, University of Waterloo, Waterloo N21 3G1, Canada and Perimeter Institute for Theoretical Physics, Waterloo, Ontario N2L 2Y5, Canada (Dated: October 30, 2017)

Machine learning vortices at the Kosterlitz-Thouless transition

Efficient and automated classification of phases from minimally processed data is one goal of machine learning in condensed matter and statistical physics. Supervised algorithms trained on raw samples of microstates can successfully detect conventional phase transitions via learning a bulk feature such as an order parameter. In this paper, we investigate whether neural networks can learn to classify phases based on topological defects. We address this question on the twodimensional classical XY model which exhibits a Kosterlitz-Thouless transition. We find significant feature engineering of the raw spin states is required to convincingly claim that features of the vortex configurations are responsible for learning the transition temperature. We further show a single-layer network does not correctly classify the phases of the XY model, while a convolutional network easily performs classification by learning the global magnetization. Finally, we design a deep network we also be of learning whoule engineering. We demonstrate the detection of

Machine Learning of Explicit Order Parameters: From the Ising Model to SU(2) Lattice Gauge Theory

Sebastian J. Wetzel¹ and Manuel Scherzer¹

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We present a procedure for reconstructing the decision function of an artificial neural network as a simple function of the input, provided the decision function is sufficiently symmetric. In this case one can easily deduce the quantity by which the neural network classifies the input. The procedure is embedded into a pipeline of machine learning algorithms able to detect the existence of different phases of matter, to determine the position of phase transitions and to find explicit expressions of the physical quantities by which the algorithm distinguishes between phases. We assume no prior knowledge about the Hamiltonian or the order parameters except Monte Carlosampled configurations. The method is applied to the Ising Model and SU(2) lattice gauge theory. In both systems we deduce the explicit expressions of the known order parameters from the decision functions of the neural networks. best classification accuracy, especially for lattices of less r systems, it remains a difficult task to learn vortices.

Machine Learning of Phase Diagrams Overview

Pro + Con -

Feed Forward Neural Network	Most powerful	Conv Layer Spatial Structure	Least Interpretable	Carrasquila, Melko, Nature 2017
Support Vector Machine	Interpretability		Not suitable for large datasets	Ponte, Melko, Phys Rev B 2017
Recurrent Neural Network	Dynamical Systems			Nieuwenburg, Bairey, Refael, Phys Rev B 2018
Principal Component Analysis	Interpretability	Most easy to use		Wang, Phys Rev B 2016
Autoencoder (Neural Network)		Conv Layer Spatial Structure		Wetzel, Phys Rev E 2017
Learning by Confusion				Nieuwenburg, Liu, Huber, Nature 2017

Supervised

Unsupervised

Hybrid

Machine Learning of Phase Diagrams Overview

New physics requires Powerful ML		s Pro +		Con - Problem	
σ	Feed Forward Neural Network	Most powerful	Conv Layer Spatial Structure	Least Interpretable	Carrasquila, Melko, Nature 2017
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нургіа	Learning by Confusion				Nieuwenburg, Liu, Huber, Nature 2017

C

Notion of Interpretability

If the neural network bases its decision on <u>one</u> single quantity/obervable Q(S)

- The larger the observable, the higher the classification probability.
- If two inputs have the same value of the observable, they have the same classification probability.
- The Neural network can be mapped via a bijective
- function to the observable

$$F(S) = f(Q(S))$$

Notion of Interpretability

- > Useful in the context of physics?
 - In Physics often only very few quantities Q(S) are characteristic features of phase transitions.
 (Renormalization Group: relevant parameters)
 - Physical Quantities are uniquely formulated by well defined formulas (in contrast to cars, faces ...)
 - Physical quantities are often highly symmetric: Rotation symmetry, translation symmetry

Interpretation of Neural Network



Wetzel, Scherzer, PRB 2017

- Interpretation Net interpolates between a general NN and a minimal optimal NN which has the same performance
- Interpretation by reducing the NN capacity in an ordered manner until one observes a performance drop
- > Inspired by extensive physical quantities (averaging layer probes for translational invariance of the quantity $\,Q(S)\,$)

Starting Neural Network:

- Conv Net with full receptive field
- > Training until converged
- Remember Loss value as measure of performance



Reinitialize the neural network with reduced receptive field sizes

- Train again until converged and compare the loss to the previous network
- Observe drop in performance from 1x2 to 1x1 and from 1x1 to baseline
- Dominant contributions must contain functions of spins and neighboring spins



Receptive Field Size	Train Loss	Validation Loss
28×28	6.1588e - 04	0.0232
1 imes 2	$1.2559\mathrm{e}\text{-}04$	1.2105 e- 07
1 imes 1	0.2015	0.1886
baseline	0.6931	0.6931

1st Network: 1x1 receptive field

Express the full neural network in 1x1 form

$$F(S) = F(\frac{1}{N}\sum_{i} f(s_i)) = \operatorname{sigmoid}(\xi(\frac{1}{N}\sum_{i} f(s_i)))$$

- > Taylor expansion eliminates all higher order terms $f(s_i) = f_0 + f_1 \ s_i + f_2 \ \underbrace{s_i^2}_1 + f_3 \ \underbrace{s_i^3}_1 + \dots$
- Regression on a single variable yields explicit form

$$F(S) \approx \text{sigmoid}\left(w \left| \frac{1}{N} \sum_{i} s_{i} \right| + b \right)$$

 \succ Where f_0, f_1 have been absorbed into weights and bias w, b

1st Network: 1x1 receptive field

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$$F(S) \approx \text{sigmoid} \left(w \left| \frac{1}{N} \sum_{i} s_i \right| + b \right)$$
 Magnetization

> Where f_0, f_1 have been absorbed into weights and bias w, b

2nd Network: 1x2 receptive field

Express the full neural network in 1x2 form

$$F(S) = F\left(\frac{1}{N}\sum_{\langle i,j \rangle_T} f(s_i, s_j)\right)$$

- > Taylor expansion contains only one addition to 1x1 case $f(s_i, s_j) = f_{0,0} + f_{1,0} s_i + f_{0,1} s_j + f_{2,0} s_i^2 + f_{1,1} s_i s_j + f_{0,2} s_j^2 + \dots$
 - Regression yields explicit form

$$D(S) \approx \text{sigmoid} \left(w \left(\frac{1}{N} \sum_{\langle i,j \rangle_T} s_i s_j \right) + b \right)$$

2nd Network: 1x2 receptive field

Express the full neural network in 1x2 form

$$F(S) = F\left(\frac{1}{N}\sum_{\langle i,j \rangle_T} f(s_i, s_j)\right)$$

- > Taylor expansion contains only one addition to 1x1 case $f(s_i, s_j) = f_{0,0} + f_{1,0} s_i + f_{0,1} s_j + f_{2,0} s_i^2 + f_{1,1} s_i s_j + f_{0,2} s_j^2 + \dots$
 - Regression yields explicit form

$$D(S) \approx \text{sigmoid} \left(w \left(\frac{1}{N} \sum_{\langle i,j \rangle_T} s_i s_j \right) + b \right) \text{Energy / 2}$$

Only half the energy since we dont sum over all neighbors

Decision functions $F(S) = \operatorname{sigmoid}(w Q(S) + b)$

$$\succ Q(S) = |1/N\sum_{i} s_i|$$

$$\Rightarrow Q(S) = \frac{1}{N} \sum_{\langle i,j \rangle_{nn}} s_i s_j$$

Deduction easily confirmed:

Perfect correlation

Note:

1x2 Network also has the Magnetization minimum which is easier to find!

Receptive Field Size	\mathbf{T}	rain Loss	Validation Loss
28×28	6.1	588e - 04	0.0232
1 imes 2	1.2	2559e-04	1.2105 e- 07
1 imes 1		0.2015	0.1886
baseline		0.6931	0.6931
Magnetization		Kashiwa, Kikuchi, Tomiya, arxiv 2019	
		Kim, Kim,	Phys Rev E 2018
			• 4

Expected Energy per site



Back to Gluons SU(2) Lattice Gauge Theory



Quarks on heavy static lattice sites.

Gluons on the connections between lattice sites are described by Matrices

 $U^x_{\mu} \in SU(2)$

SU(2) Lattice Gauge Theory

Describes smallest loop on the lattice

$$S_{\text{Wilson}}[U] = \beta_{\text{latt}} \sum_{x} \sum_{\mu < \nu} \text{Re tr} \left(1 - U_{\mu\nu}^x\right)$$

$$U^{x}_{\mu\nu} = U^{x}_{\mu}U^{x+\hat{\mu}}_{\nu}U^{x+\hat{\mu}+\hat{\nu}}_{-\mu}U^{x+\hat{\nu}}_{-\nu}$$

$$U^x_{\mu} \in SU(2)$$

$$U_{\mu}^{x} = a_{\mu}^{x} 1 + i \left(b_{\mu}^{x} \sigma_{1} + c_{\mu}^{x} \sigma_{2} + d_{\mu}^{x} \sigma_{3} \right)$$

 Each Matrix is parametrized by 4 real numbers.

We performed a MC simulation on a lattice of size 8x8x8x2 as input for the Neural Network Each Matrix connects two lattice sites

- Toy model for confinement of particles in atomic cores.
- Polyakov Loop is Order
 Parameter for in the limit of infinitely heavy quarks.
 - Perfect Testing Ground: Polyakov Loop Order Parameter is non-linear and non-local.

Unsupervised Learning (PCA) SU(2) Lattice Gauge Theory

$$S_{\text{Wilson}}[U] = \beta_{\text{latt}} \sum_{x} \sum_{\mu < \nu} \text{Re tr} \left(1 - U_{\mu\nu}^x \right)$$

- Latent parameter does not correspond to order parameter
 - PCA + Reconstruction loss can be used to infer different phases

Training at phase indications from unsupervised learning (wait for next slide)



SU(2) Lattice Gauge Theory

$$S_{\text{Wilson}}[U] = \beta_{\text{latt}} \sum_{x} \sum_{\mu < \nu} \text{Re tr} \left(1 - U_{\mu\nu}^x \right)$$

 Find phase transition close to lattice calculation

Prediction is inaccurate: Monte Caro Simulations not thermalized

Training at phase indications from unsupervised learning

Testing in interval containing phase transition



Interpretation of Neural Network SU(2) Gauge Theory (2x8x8x8 Lattice)

General decision function:	
$F(S) = \operatorname{sigmoid}(w Q(S) + b)$)

2x1x1x1 Decision function:

Receptive Field Size	Train Loss	Validation Loss
$2 \times 8 \times 8 \times 8$	1.0004e - 04	2.6266e - 04
2 imes 1 imes 1 imes 1	8.8104 e-08	6.8276e-08
$2 imes 1 imes 1 imes 1^{*}$	2.2292 e- 07	$4.2958\mathrm{e}{\textbf{-}07}$
$1 \times 1 \times 1 \times 1$	0.6620	0.9482
baseline	0.6931	0.6931

$$F(S) \approx \text{sigmoid}\left(w\left(\frac{2}{N}\sum_{\vec{x}}f(\{U_{\mu}^{x_{0},\vec{x}}\})\right) + b\right)$$

Regression yields 561 terms: $f(\{U^{x_0}_{\mu}\}) \approx +7.3816 \ a^0_{\tau}a^1_{\tau} + 0.2529 \ a^1_{\tau}b^1_{\tau} + \dots$ $-0.2869 \ d^0_{\tau}c^1_{\tau} - 7.2279 \ b^0_{\tau}b^1_{\tau}$ $-7.3005 \ c^0_{\tau}c^1_{\tau} - 7.4642 \ d^0_{\tau}d^1_{\tau}$

$$f(\{U^{x_0}_{\mu}\}) = a^0_{\tau} a^1_{\tau} - b^0_{\tau} b^1_{\tau} - c^0_{\tau} c^1_{\tau} - d^0_{\tau} d^1_{\tau} = \operatorname{tr}\left(U^0_{\tau} U^1_{\tau}\right)$$

Neural Network uses Polyakov Loop to distinguish between phases.

Interpretation of Neural Network SU(2) Gauge Theory (2x8x8x8 Lattice)



Polyakov Loop

Interpretation of Neural Network SU(2) Gauge Theory (2x8x8x8 Lattice)



Polyakov Loop

Note: We have constructed the PL without prior knowledge!

Conclusion

Neural Networks are capable of producing phase diagrams for many physical systems.

- NNs are no longer a black box algorithm in the context of order parameter based phase transitions.
- Neural Networks learn the same physical quantities that we humans use (Landau/Ehrenfest)
 - In the spirit of the conference: robust features
- In some cases we can determine the nature of phases by constructively interpreting what neural networks learn.